# MAT 520: Problem Set 10

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Provide an example for a non-normal operator  $A \in \mathcal{B}(\mathcal{H})$  and a point in the resolvent set  $z \in \rho(A)$  where

$$
||(A - z\mathbb{1})^{-1}|| \le \frac{1}{\text{dist}(z, \sigma(A))}
$$

does not hold.

#### Solution

**Proof.** Let  $\mathcal{H} = \mathbb{R}^2$  and  $A \in \mathcal{B}(\mathcal{H})$  be the operator given by the matrix

$$
A:=\begin{bmatrix}2&1\\0&2\end{bmatrix}
$$

Then, we know that  $\sigma(A) = \{2\}$ . So, for  $z = 1 \in \rho(A)$ , we see that

$$
(A - 1)^{-1} = \begin{bmatrix} 1 & 1 \\ 0 & 1 \end{bmatrix}^{-1} = \begin{bmatrix} 1 & -1 \\ 0 & 1 \end{bmatrix}
$$

We may compute the operator norm using the fact that for an  $n \times n$  matrix B, we have  $||B||_{op} = \sqrt{\lambda_{max}(B^*B)}$ is the square root of the largest eigenvalue of  $B^*B$ . So, letting  $B := (A - 1)^{-1}$ , we seek eigenvalues of the matrix

$$
B^*B = \begin{bmatrix} 1 & 0 \\ -1 & 1 \end{bmatrix} \begin{bmatrix} 1 & -1 \\ 0 & 1 \end{bmatrix} = \begin{bmatrix} 1 & -1 \\ -1 & 2 \end{bmatrix}
$$

We see that  $\lambda$  is an eigenvalue iff  $(1 - \lambda)(2 - \lambda) - 1 = 0 \iff \lambda^2 - 3\lambda + 1 = 0 \iff \lambda = \frac{3 \pm \sqrt{5}}{2}$ . Therefore,  $\lambda_{\text{max}}(B^*B) = \frac{3+\sqrt{5}}{2}$ , and so

$$
\|(A - 1)^{-1}\| = \sqrt{\frac{3 + \sqrt{5}}{2}} \approx 1.618 > 1
$$

However, we have that

$$
\frac{1}{\text{dist}(z, \sigma(A))} = \frac{1}{|1 - 2|} = 1
$$

Therefore, for this choice of  $H$ ,  $A$ , and  $z$ , we have that

$$
\|(A - z\mathbb{1})^{-1}\| > \frac{1}{\text{dist}(z, \sigma(A))}
$$

 $\blacksquare$ 

Let  $A \in \mathcal{B}(\mathcal{H})$  be an operator with polar decomposition  $A = U|A|$ . Define functions  $f_n : [0, \infty) \to \mathbb{R}$  via

$$
f_n(x) := \begin{cases} \frac{1}{x} & x \ge \frac{1}{n} \\ n & x \le \frac{1}{n} \end{cases}
$$

Prove that

$$
U = \operatorname*{s-lim}_{n \to \infty} Af_n(|A|)
$$

### Solution

Proof. We wish to show that

$$
1 = \mathop{\rm s-lim}_{n \to \infty} |A| f_n(|A|),
$$

since the main result will follow by applying U to both sides. Since  $|A|$  is self-adjoint, we may apply the functional calculus. Define  $g_n : [0, \infty) \to \mathbb{R}$  via

$$
g_n(x) := x f_n(x) = \begin{cases} 1 & x \ge \frac{1}{n} \\ nx & x \le \frac{1}{n} \end{cases}
$$

Letting  $I : [0, \infty) \to [0, \infty)$  be the identity map  $x \mapsto x$ , we have that  $g_n(x) = I(x) f_n(x)$ . Then, by the homomorphism property of the functional calculus and the fact that  $I(|A|) = |A|$ ,

$$
g_n(|A|) = I(|A|)f_n(|A|) = |A|f_n(|A|)
$$

So, we want to show that s-lim<sub>n→∞</sub>  $g_n(|A|) = \mathbb{1}$ , and the result will follow. To see this, we simply note that  $g_n(x)$  converges to 1 for every  $x \in (0,\infty)$ , and so  $g_n \to 1$  pointwise a.e. on  $[0,\infty)$ . Then, over  $\sigma(|A|)$  we see that  $g_n \to 1$  pointwise and  $||g_n||_{\infty} \leq 1 < \infty$ . So, the measurable functional calculus (Theorem VII.2(d) in Reed & Simon) gives that  $g_n(|A|) \to \mathbb{1}$  strongly. The main result follows.

Prove that if  $A \in \mathcal{B}(\mathcal{H})$  is normal, then

 $||A|| = r(A)$ 

where  $r(\cdot)$  is the spectral radius.

### Solution

**Proof.** By the spectral mapping theorem and the continuous functional calculus on the function  $z \mapsto |z|^2$ (which we may apply since A is normal, see Theorem 8.40 in the lecture notes), we have that

$$
\sigma(|A|^2)=\{|z|^2:\;z\in\sigma(A)\}
$$

Thus,

$$
r(|A|^2) = \sup_{z \in \sigma(A)} \{|z|^2\},\,
$$

and so

$$
r(A) = \sup_{z \in \sigma(A)} \{|z|\} = \sqrt{r(|A|^2)}
$$

We know that  $||A||^2 = r(|A|^2)$  by Theorem 8.6 in the lecture notes since  $|A|^2$  is self-adjoint. So, by this and the  $C^*$  identity,

$$
r(A) = \sqrt{\|A\|^2\|} = \sqrt{\|A\|^2} = \|A\|
$$

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Let  $A \in \mathcal{B}(\mathcal{H})$  be normal. Show there exists some finite measure space  $(M, \mu)$  and a unitary  $U : \mathcal{H} \to$  $L^2(M,\mu)$  such that there exists a bounded Borel function  $f: M \to \mathbb{C}$  such that

$$
(UAU^*\psi)(m) = f(m)\psi(m) \quad (m \in M, \ \psi \in L^2(M, \mu))
$$

### Solution

Proof. We may express A as a linear combination of two self-adjoint operators via

 $A = \mathbb{Re}{A} + i \mathbb{Im}{A}$ 

Furthermore, since A is normal we see that these two operators commute. Applying the result of Problem 5 with  $A_1 := \mathbb{Re}{A}$  and  $A_2 := \mathbb{Im}{A}$  yields the finite measure space and unitary such that  $\mathbb{Re}{A}$  is represented as multiplication by f and  $\text{Im } \{A\}$  is represented as multiplication by g. A is therefore represented as multiplication by  $f + ig : M \to \mathbb{C}$  as desired.

Show that if  $A, B \in \mathcal{B}(\mathcal{H})$  are two self-adjoint operators such that  $[A, B] = 0$ , then there exists a finite measure space  $(M, \mu)$  and a unitary  $U : \mathcal{H} \to L^2(M, \mu)$  such that there are two bounded Borel functions  $f, g: M \to \mathbb{R}$  which obey

$$
(UAU^*\psi)(m) = f(m)\psi(m)
$$
  

$$
(UBU^*\psi)(m) = g(m)\psi(m)
$$

for all  $m \in M$  and  $\psi \in L^2(M, \mu)$ .

#### Solution

**Proof.** Suppose that A and B commute. We claim that for all  $\Omega \subseteq \sigma(A)$  and  $\Sigma \subseteq \sigma(B)$  Borel it holds that  $[\chi_{\Omega}(A), \chi_{\Sigma}(B)] = 0$ , or in other words that the spectral projections commute. Note that for each continuous function  $f : \sigma(A) \to \mathbb{C}$ , we have that

$$
f(A)B = Bf(A)
$$

since we may uniformly approximate  $f$  by polynomials (Stone-Weierstrass) and the above trivially holds for polynomial functions. For any  $\psi \in \mathcal{H}$  we may approximate  $\chi_{\Omega}(\cdot)$  in the  $\mu_{A,\psi}$  measure by continuous functions on  $\sigma(A)$ , which reveals that

$$
\chi_{\Omega}(A)B\psi = B\chi_{\Omega}(A)\psi
$$

since the region of error from the approximation can be made to have approximately small  $\mu_{A,\psi}$  measure. Since this holds for all  $\psi \in \mathcal{H}$ , we see that

$$
\chi_{\Omega}(A)B = B\chi_{\Omega}(A)
$$

which gives that

$$
\chi_{\Omega}(A)\chi_{\Sigma}(B) = \chi_{\Sigma}(B)\chi_{\Omega}(A)
$$

for all Borel  $\Omega \subseteq \sigma(A)$  and  $\Sigma \subseteq \sigma(B)$ . Define the set of rectangles

$$
\mathcal{R} := \{ E \subseteq \mathbb{R}^2 : E = \Omega \times \Sigma \text{ for } \Omega \subseteq \sigma(A) \text{ and } \Sigma \subseteq \sigma(B) \text{ Borel} \}
$$

We will define a functional calculus starting with simple functions that can be written as linear combinations of characteristic functions of disjoint rectangles from  $R$ . Write

$$
\mathcal{S} := \left\{ f : f = \sum_{i=1}^{n} \alpha_i \chi_{R_i} \text{ with } \{R_i\}_i \subseteq \mathcal{R} \text{ pairwise disjoint} \right\}
$$

to be the set of simple functions. For such  $f \in \mathcal{S}$ , we define

$$
f(A, B) := \sum_{i=1}^{n} \alpha_i \chi_{\Omega_i}(A) \chi_{\Sigma_i}(B)
$$

where  $R_i = \Omega_i \times \Sigma_i$ . We stop to note that for all such f,  $||f(A, B)|| \le ||f||_{\infty}$ . To see this, observe that  $||f||_{\infty} = \max_i {|\alpha_i|}$ . Since the  $\Omega_i$ 's are disjoint from each other and similarly for the  $\Sigma_i$ 's, we see that the projection operators  $\{\chi_{\Omega_i}(A)\chi_{\Sigma_i}(B)\}_i$  are all pairwise orthogonal. So, for any  $\psi \in \mathcal{H}$  we that

$$
||f(A,B)\psi||^2 = \sum_{i=1}^n |\alpha_i|^2 ||\chi_{\Omega_i}(A)\chi_{\Sigma_i}(B)\psi||^2 \le ||f||^2_{\infty} \sum_{i=1}^n ||\chi_{\Omega_i}(A)\chi_{\Sigma_i}(B)\psi||^2 \le ||f||^2_{\infty} ||\psi||^2,
$$

where the last inequality comes from the Pythagorean theorem. So, this defines a functional calculus

$$
\phi : \mathcal{S} \to \mathcal{B}(\mathcal{H})
$$

Problem 5 continued on next page... 6

which is clearly a linear transformation and satisfies  $\|\phi(f)\| \leq \|f\|_{\infty}$ . Thus, since the set S is dense in  $C(\sigma(A) \times \sigma(B))$  equipped with the uniform norm (we may apply Stone-Weierstrass since S separates points), we may use the BLT theorem to construct a continuous functional calculus  $\phi : C(\sigma(A) \times \sigma(B)) \to \mathcal{B}(\mathcal{H})$ .

We now construct a Borelian functional calculus. For any  $\psi \in \mathcal{H}$  we see that the map

$$
C(\sigma(A) \times \sigma(B)) \ni f \mapsto \langle \psi, \phi(f)\psi \rangle = \langle \psi, f(A, B)\psi \rangle
$$

is a real-valued, continuous linear functional on  $C(\sigma(A) \times \sigma(B))$ . Therefore, by the Riesz-Markov theorem, there is a unique Borel measure  $\mu_{\psi}$  on  $\mathbb{R}^2$  with  $\mu_{\psi}(\mathbb{R}^2) = ||\psi||^2$  and

$$
\langle \psi, f(A,B) \psi \rangle = \int_{\mathbb{R}^2} f(z) d\mu_{\psi}(z)
$$

These measures  $\mu_{\psi}$  are spectral measures, and we may use the polarization identity to uniquely define  $\langle \psi, f(A, B)\varphi \rangle$  for any bounded, Borel-measurable function  $f : \sigma(A) \times \sigma(B) \to \mathbb{R}$ . Thus upgrades us to a Borelian functional calculus.

We proceed to a spectral theorem. Call a vector  $\psi \in \mathcal{H}$  cyclic for  $(A, B)$  if span $\{f(A, B)\psi : f \in$  $C(\sigma(A) \times \sigma(B))$  is dense in H. Then, by basically the same proof as Lemma 10.22 in the lecture notes, we see that if  $\psi$  is cyclic for  $(A, B)$  then there is a unitary operator  $U: \mathcal{H} \to L^2(\sigma(A) \times \sigma(B), \mu_{\psi})$  for which

$$
(UAU^* f)(x_1, x_2) = x_1 f(x_1, x_2) \quad ((x_1, x_2) \in \sigma(A) \times \sigma(B))
$$
  

$$
(UBU^* f)(x_1, x_2) = x_2 f(x_1, x_2) \quad ((x_1, x_2) \in \sigma(A) \times \sigma(B))
$$

Now, all we must do is decompose  $H$  into a direct sum of spaces which each admit a cyclic vector, which can be done cleanly via Zorn's lemma. So, we see that

$$
\mathcal{H}=\bigoplus_{n=1}^N\mathcal{H}_n
$$

where  $N \in \mathbb{N} \cup \{\infty\}$  and each  $\mathcal{H}_n$  admits a cyclic vector  $\psi_n$ . We see that H is unitarily equivalent to  $\bigoplus_{n=1}^N L^2(\sigma(A) \times \sigma(B), \mu_{\psi_n})$  and for all  $n = 1, ..., N$  and  $f_n \in L^2(\sigma(A) \times \sigma(B), \mu_{\psi_n}),$ 

$$
(UAU^* f)_n(x_1, x_2) = x_1 f_n(x_1, x_2) \quad ((x_1, x_2) \in \sigma(A) \times \sigma(B))
$$
  

$$
(UBU^* f)_n(x_1, x_2) = x_2 f_n(x_1, x_2) \quad ((x_1, x_2) \in \sigma(A) \times \sigma(B))
$$

Set  $M := \bigoplus_{n=1}^{N} (\sigma(A) \times \sigma(B))$  and  $\mu := \bigotimes_{n=1}^{N} \mu_{\psi_n}$ , The above statment now reads that, for every  $f \in$  $L^2(M,\mu)$ , we have

$$
(UAU^*f)(m) = F_A(m)f(m) \quad (m \in M)
$$
  

$$
(UBU^*f)(m) = F_B(m)f(m) \quad (m \in M)
$$

where  $F_A$  acts on the  $\sigma(A)$  part of each constituent  $\sigma(A) \times \sigma(B)$  in M and  $F_B$  acts on the  $\sigma(B)$  parts. We note that  $\mu(M) = \sum_{n=1}^{N} ||\psi_n||^2$ ; if we select  $\psi_n$ 's such that  $||\psi_n|| < 2^{-n}$  (which we may do since scaling does not change cyclicity in  $\mathcal{H}_n$ , we see that  $\mu(M) \leq 1$ .

Prove that for  $A \in \mathcal{B}(\mathcal{H})$  self-adjoint and  $\chi(A)$  the projection-valued measure of A, we have

 $\lambda \in \sigma(A) \iff \chi_{(\lambda-\varepsilon,\lambda+\varepsilon)}(A) \neq 0 \quad (\varepsilon > 0)$ 

#### Solution

**Proof.** We will show the contrapositives for both directions. Namely, we prove the following:

$$
\lambda \in \rho(A) \iff \exists \varepsilon > 0 \text{ s.t. } \chi_{(\lambda - \varepsilon, \lambda + \varepsilon)}(A) = 0
$$

 $(\iff)$  Suppose that  $\chi_{(\lambda-\varepsilon,\lambda+\varepsilon)}(A)=0$  for some  $\varepsilon>0$ ; we want to show that this implies  $\lambda\in\rho(A)$ . Note that for all  $\varphi, \psi \in \mathcal{H}$ ,

$$
0 = \langle \varphi, \chi_{(\lambda - \varepsilon, \lambda + \varepsilon)}(A)\psi \rangle = \mu_{A, \varphi, \psi}((\lambda - \varepsilon, \lambda + \varepsilon))
$$

Therefore, for all  $\varphi, \psi \in \mathcal{H}$  we see that for all  $z \in \text{spt}(\mu_{A,\varphi,\psi})$  we have  $|z - \lambda| \geq \varepsilon$ , and so the map sending  $z \mapsto \frac{1}{z-\lambda}$  is bounded on  $\text{spt}(\mu_{A,\varphi,\psi})$ . In other words, we have that  $\left|\frac{1}{z-\lambda}\right| \leq \frac{1}{\varepsilon}$  for  $\mu_{A,\varphi,\psi}$ -a.e.  $z \in \mathbb{R}$ . By the bounded measurable functional calculus, we therefore see that

$$
\langle \varphi, (A - \lambda \mathbb{1})^{-1} \psi \rangle = \int_{\mathbb{R}} \frac{1}{z - \lambda} d\mu_{A, \varphi, \psi}(z)
$$

converges for all  $\varphi, \psi \in \mathcal{H}$ . Since all the matrix elements are defined, the resolvent  $(A - \lambda \mathbb{1})^{-1}$  exists. In particular,  $\lambda \in \rho(A)$ .

 $(\implies)$  Suppose that  $\lambda \in \rho(A)$ . Since  $\rho(A)$  is open, there is an  $\varepsilon > 0$  such that  $(\lambda - \varepsilon, \lambda + \varepsilon) \subseteq \rho(A)$ . Since  $\text{spt}(\mu_{A,\psi}) \subseteq \sigma(A)$  for all  $\psi \in \mathcal{H}$ , this reveals that for all  $\psi \in \mathcal{H}$ ,

$$
(\lambda - \varepsilon, \lambda + \varepsilon) \subseteq \text{spt}(\mu_{A,\psi})^C \implies \int_{\mathbb{R}} \chi_{(\lambda - \varepsilon, \lambda + \varepsilon)}(z) d\mu_{A,\psi}(z) = 0
$$

By construction of the measurable functional calculus, for all  $\psi \in \mathcal{H}$  we have

$$
\langle \psi, \chi_{(\lambda-\varepsilon,\lambda+\varepsilon)}(A)\psi \rangle = \int_{\mathbb{R}} \chi_{(\lambda-\varepsilon,\lambda+\varepsilon)}(z) d\mu_{A,\psi}(z) = 0
$$

So, by Theorem 7.11 in the lecture notes, since all the diagonal elements of  $\chi_{(\lambda-\varepsilon,\lambda+\varepsilon)}(A)$  are 0 we know that  $\chi_{(\lambda-\varepsilon,\lambda+\varepsilon)}(A) = 0$ . So, there exists an  $\varepsilon > 0$  such that  $\chi_{(\lambda-\varepsilon,\lambda+\varepsilon)}(A) = 0$ .

Prove that the **only** norm-closed  $\ast$ -ideals in  $\mathcal{B}(\mathcal{H})$  are  $\{0\}$ ,  $\mathcal{K}(\mathcal{H})$ , and  $\mathcal{B}(\mathcal{H})$ .

#### Solution

**Proof.** Let  $A \subseteq \mathcal{B}(\mathcal{H})$  be a norm-closed \*-ideal. Clearly,  $0 \in \mathcal{A}$ . Also,  $\mathcal{A}$  is a vector subspace. We claim the following:

- (1) If A contains any nonzero operator, then  $\mathcal{K}(\mathcal{H}) \subseteq \mathcal{A}$
- (2) If A contains any noncompact operator, then  $A = \mathcal{B}(\mathcal{H})$

The result follows from these facts. To see this, we will exhaust all the possible cases. Suppose first that  $A$ contains no nonzero compact operator and no nonzero noncompact operator; then,  $\mathcal{A} = \{0\}$ . Next, suppose that A contains some nonzero compact operator (and so  $\mathcal{K}(\mathcal{H}) \subseteq \mathcal{A}$  by (1)), but no nonzero noncompact operator. Then,  $\mathcal{B}(\mathcal{H}) \setminus \mathcal{K}(\mathcal{H}) \subseteq \mathcal{A}^C \implies \mathcal{A} \subseteq \mathcal{K}(\mathcal{H})$ , and so  $\mathcal{A} = \mathcal{K}(\mathcal{H})$ . Lastly, suppose that A contains some noncompact operator; (2) then implies that  $\mathcal{A} = \mathcal{B}(\mathcal{H})$ . So, to complete the proof it suffices to show (1) and (2).

(1) We will show that A contains every rank-one operator of the form  $\varphi \otimes \psi^*, \varphi, \psi \in \mathcal{H}$ , since then A will contain every finite-rank operator by properties of a vector space. From this, we will see that  $A$  contains  $\mathcal{K}(\mathcal{H})$  by the fact that  $\mathcal{K}(\mathcal{H})$  is the norm-closure of the set of finite-rank operators and A is norm-closed. So, let  $\varphi, \psi \in \mathcal{H}$  be arbitrary. Let  $B \in \mathcal{A}$  be nonzero. Then, there is some  $\phi \in \mathcal{H}$  such that  $B(\phi)$  is nonzero. By the Hahn-Banach theorem (linear functionals separate points), there is some  $\eta \in \mathcal{H}$  such that  $\langle \eta, B(\phi) \rangle = 1$ . We claim that

$$
\varphi \otimes \psi^* = (\varphi \otimes \eta^*) B(\phi \otimes \psi^*),
$$

which by the two-sided-ideal property would mean that  $\varphi \otimes \psi^* \in \mathcal{A}$ . To see this, note that for any  $\xi \in \mathcal{H}$ ,

$$
(\varphi \otimes \eta^*)B(\phi \otimes \psi^*)(\xi) = (\varphi \otimes \eta^*)B(\langle \psi, \xi \rangle \phi) = \langle \psi, \xi \rangle \langle \eta, B(\phi) \rangle \varphi = \langle \psi, \xi \rangle \varphi = (\varphi \otimes \psi^*)(\xi)
$$

So,  $\varphi \otimes \psi^* \in \mathcal{A}$  for all  $\varphi, \psi \in \mathcal{H}$ . Thus,  $\mathcal{A}$  contains all the finite-rank operators, and by norm-closure we have that  $\mathcal{K}(\mathcal{H}) \subseteq \mathcal{A}$ .

(2) Suppose now that A contains a noncompact operator  $S \in \mathcal{A}$ . By noncompactness, im(S) must have a closed infinite-dimensional subspace, call it M. Define the closed vector subspace  $N := S^{-1}(M) \cap \text{ker}(S)^{\perp}$ . Then, consider the restricted bounded linear operator  $S|_N : N \to M$ . By construction,  $S|_N$  is both injective (since ker(S) ∩ N = {0}) and surjective (since im(S|<sub>N</sub>) = M). Since N and M are both closed subspaces of H they are Banach spaces, and so  $S|_N^{-1}: M \to N$  is continuous by the inverse mapping theorem. Let  $P_M$ be the orthogonal projection onto M, and define the map  $T \in \mathcal{B}(\mathcal{H})$  via

$$
T = S|_N^{-1} P_M
$$

Then, for all  $\varphi \in M$  we see that  $ST\varphi = \varphi$  whereas for all  $\varphi \in M^{\perp}$  we see that  $ST\varphi = 0$ . So,  $ST = P_M$ . Now, let  $R \in \mathcal{B}(\mathcal{H})$  be arbitrary; we will show that  $R \in \mathcal{A}$ . Let  $\{\varphi_n\}_{n \in \mathbb{N}} \subseteq M$  be an orthonormal basis of M and let  $\{\psi_n\}_{n=1}^N \subseteq \overline{\text{im}(R)}$  be an orthonormal basis of  $\overline{\text{im}(R)}$  (N may be finite or countably infinite, which is fine since  $\{\varphi_n\}_n$  is infinite). Let  $U \in \mathcal{B}(\mathcal{H})$  be the map sending  $\psi_n \mapsto \varphi_n$  and extended linearly, such that  $U \equiv 0$  over  $\overline{\text{im}(R)}^{\perp}$ . Then, U is a partial isometry and  $\text{im}(U) \subseteq M$ . Note that for all  $\xi \in \mathcal{H}$ , we have that

$$
R\xi \in \ker(U)^{\perp} \implies R\xi = |U|^2 R\xi = U^* U R\xi
$$

However, we know that  $UR\xi \in M$  and so  $STUR\xi = UR\xi$ . Thus,  $R\xi = U^*STUR\xi$ . Since this holds for all  $\xi \in \mathcal{H}$ , we see that  $R = U^*STUR$ , which means that  $R \in \mathcal{A}$  by the two-sided-ideal property. Thus,  $\mathcal{A} = \mathcal{B}(\mathcal{H})$ .

Let  $\mathcal{H} := \ell^2(\mathbb{Z})$  and on it define the discrete Laplacian

 $-\Delta := 21 - R - R^*$ 

where  $R$  is the bilateral right shift operator.

- (a) For  $x \in \mathbb{Z}$ , is  $\delta_x$  a cyclic vector for  $-\Delta$ ?
- (b) Define  $f: \mathbb{C}^+ \to \mathbb{C}$  via

$$
f(z) = \langle \delta_0, (-\Delta - z\mathbb{1})^{-1} \delta_0 \rangle
$$

Find an explicit expression for f using the Fourier series.

(c) Calculate

$$
\lim_{\varepsilon\to 0^+}\mathbb{I}\mathrm{m}\left\{f(E+i\varepsilon)\right\}
$$

for the two cases  $E \in (0, 4)$  and  $E \in \mathbb{R} \setminus (0, 4)$ .

(d) Calculate the spectral measure of  $(-\Delta, \delta_0)$  and determine its type (with respect to the Lebesgue decomposition theorem where the reference measure is the Lebesgue measure, i.e., ac, sc, or pp).

#### Solution

**Proof.** (a) Pick an  $x \in \mathbb{Z}$ . Then, we see that applying  $-\Delta$  symmetrically introduces positions to the left and right of x. Put differently, we have that for all  $n, k \in \mathbb{N}$ ,

$$
\langle \delta_{x+k}, (-\Delta)^n \delta_x \rangle = \langle \delta_{x-k}, (-\Delta)^n \delta_x \rangle
$$

By linearity and continuity of the inner product, we see that for all  $\psi \in \overline{\text{span}\{(-\Delta)^n \delta_x : n \in \mathbb{N}\}}$  and all  $k \in \mathbb{N}$  we have

$$
\langle \delta_{x+k}, \psi \rangle = \langle \delta_{x-k}, \psi \rangle
$$

Therefore we clearly cannot have that  $\delta_{x+k} \in \overline{\text{span}\{(-\Delta)^n \delta_x : n \in \mathbb{N}\}}$  for any  $k \in \mathbb{N}$ , which means that  $\overline{\text{span}\{(-\Delta)^n \delta_x : n \in \mathbb{N}\}}$  is not dense. Thus,  $\delta_x$  isn't cyclic.

(b) Let  $\mathcal{F}: \mathcal{H} \to L^2(\mathbb{S}^1)$  be the Fourier transform from Problem Set 8. Then, since it is unitary we have that for all  $z \in \mathbb{C}^+$ ,

$$
f(z) = \langle \delta_0, (-\Delta - z\mathbb{1})^{-1} \delta_0 \rangle = \langle \mathcal{F}\delta_0, \mathcal{F}(-\Delta - z\mathbb{1})^{-1} \mathcal{F}^* \mathcal{F}\delta_0 \rangle_{L^2(\mathbb{S}^1)}
$$

We know that  $(\mathcal{F}\delta_0)(x) = e^{-i0x} = 1$ , the constant 1 function. Also, we know from Problem Set 8 and linearity that  $\mathcal{F}(-\Delta)\mathcal{F}^*$  is the multiplication operator by the function  $\theta \mapsto 2-2\cos(\theta)$ . So, by the functional calculus we know that  $\mathcal{F}(-\Delta - z\mathbb{1})^{-1}\mathcal{F}^*$  is the multiplication operator by the function  $\theta \mapsto \frac{1}{2-2\cos(\theta)-z}$ . So, we may compute this inner product in  $L^2(\mathbb{S}^1)$  to see that for  $z \in \mathbb{C}^+,$ 

$$
f(z) = \frac{1}{2\pi} \int_0^{2\pi} \overline{(\mathcal{F}\delta_0)(\theta)} \frac{1}{2 - 2\cos(\theta) - z} (\mathcal{F}\delta_0)(\theta) d\theta
$$
  
= 
$$
\frac{1}{2\pi} \int_0^{2\pi} \frac{1}{(2 - z) - 2\cos(\theta)} d\theta
$$

It is a well-known fact from residue integration that  $\int_0^{2\pi} \frac{1}{a+b\cos(\theta)} d\theta = \frac{2\pi}{\sqrt{a^2-b^2}}$ , where the result may be complex for general  $a, b \in \mathbb{C}$ . Therefore, we have that

$$
f(z) = \frac{1}{\sqrt{(2-z)^2 - 4}} = \frac{1}{\sqrt{z^2 - 4z}} \quad (z \in \mathbb{C}^+)
$$

Problem 8 continued on next page... 10

(c) Let  $E \in \mathbb{R}$  and  $\varepsilon > 0$ . Noting that  $\mathbb{I}$ m  $\{\cdot\}, \sqrt{\cdot}$ , and  $\frac{1}{\cdot}$  are all continuous away from 0, we see that when  $E \notin \{0, 4\}$  we may directly compute the limit as

$$
\lim_{\varepsilon \to 0^+} \text{Im} \left\{ f(E + i\varepsilon) \right\} = \text{Im} \left\{ \frac{1}{\sqrt{\lim_{\varepsilon \to 0^+} (E + i\varepsilon)^2 - 4(E + i\varepsilon)}} \right\} = \text{Im} \left\{ \frac{1}{\sqrt{E^2 - 4E}} \right\}
$$

In particular, when  $E > 4$  or  $E < 0$  the imaginary part is 0. When  $E \in (0, 4)$  we have that

$$
\operatorname{Im}\left\{\frac{1}{\sqrt{E^2 - 4E}}\right\} = \operatorname{Im}\left\{\frac{1}{i\sqrt{4E - E^2}}\right\} = \frac{1}{\sqrt{4E - E^2}} > 0
$$

When E is 0 or 4 we see that  $(E + i\varepsilon)^2 - 4(E + i\varepsilon) = E^2 - \varepsilon^2 + 2Ei\varepsilon - 4E - 4i\varepsilon = -\varepsilon^2 \pm 4i\varepsilon$ , and so

$$
\operatorname{Im}\left\{f(E+i\varepsilon)\right\} = \operatorname{Im}\left\{\frac{1}{\sqrt{-\varepsilon^2 \pm 4i\varepsilon}}\right\} = \operatorname{Im}\left\{\sqrt{\frac{1}{-\varepsilon^2 \pm 4i\varepsilon}}\right\}
$$

Clearly, as  $\varepsilon \to 0^+$  we see that this approaches infinity. To summarize,

$$
\lim_{\varepsilon \to 0^+} \text{Im} \left\{ f(E + i\varepsilon) \right\} = \begin{cases} 0 & E \in \mathbb{R} \setminus [0, 4] \\ \frac{1}{\sqrt{4E - E^2}} & E \in (0, 4) \\ \infty & E \in \{0, 4\} \end{cases}
$$

We take the time here to also note that for  $E \in \{0, 4\}$ ,

$$
\lim_{\varepsilon \to 0^+} \varepsilon \operatorname{Im} \left\{ f(E + i\varepsilon) \right\} = \lim_{\varepsilon \to 0^+} \operatorname{Im} \left\{ \sqrt{\frac{1}{-1 \pm \frac{4i}{\varepsilon}}} \right\} = 0
$$

(d) Let  $\mu_{-\Delta,\delta_0}$  be the spectral measure of  $(-\Delta,\delta_0)$ . Write the Radon-Nikodym decomposition of  $\mu_{-\Delta,\delta_0}$ w.r.t. the Lebesgue measure as

$$
\mu_{-\Delta,\delta_0} = f\lambda + \mu_{\text{sing}}
$$

By Lemma 10.10 in the lecture notes and part (c), we see that

$$
f(\lambda) = \lim_{\varepsilon \to 0^+} \frac{1}{\pi} \operatorname{Im} \{ f(\lambda + i\varepsilon) \} = \begin{cases} 0 & \lambda \in \mathbb{R} \setminus [0, 4] \\ \frac{1}{\pi \sqrt{4\lambda - \lambda^2}} & \lambda \in (0, 4) \\ \infty & \lambda \in \{0, 4\} \end{cases}
$$

By Lemma 10.11 in the lecture notes,  $spt(\mu_{sing}) = \{0, 4\}$ , but at these values  $\mu_{-\Delta, \delta_0}$  has no point masses. Thus, the spectral measure  $\mu_{-\Delta,\delta_0}$  is absolutely continuous w.r.t. Lebesgue with support on [0,4] and a density given by  $f$  above.  $\blacksquare$ 

Let  $\mathcal{H} := \ell^2(\mathbb{Z})$  and on it define the multiplication operator  $V(X)$  via

$$
(V(X)\psi)(x) := V(x)\psi(x) \quad (x \in \mathbb{Z}, \psi \in \mathcal{H}),
$$

where  $V : \mathbb{Z} \to \mathbb{R}$  is some bounded sequence.

- (a) For  $x \in \mathbb{Z}$ , is  $\delta_x$  a cyclic vector for  $V(X)$ ?
- (b) For any  $x \in \mathbb{Z}$ , define  $f_x : \mathbb{C}^+ \to \mathbb{C}$  via

$$
f_x(z) = \langle \delta_x, (V(X) - z\mathbb{1})^{-1} \delta_x \rangle
$$

Find an explicit expression for  $f_x$ .

(c) Calculate both

$$
\lim_{\varepsilon \to 0^+} \text{Im} \left\{ f_x(E + i\varepsilon) \right\}
$$

and

$$
\lim_{\varepsilon \to 0^+} \varepsilon \operatorname{Im} \left\{ f_x(E + i\varepsilon) \right\}
$$

for all  $E \in \mathbb{R}$  (separate into cases).

(d) Calculate the spectral measure of  $(V(X), \delta_0)$  and determine its type (with respect to the Lebesgue decomposition theorem where the reference measure is the Lebesgue measure, i.e., ac, sc, or pp).

#### Solution

**Proof.** (a) For any  $x \in \mathbb{Z}$ , note that  $V(X)\delta_x = V(x)\delta_x$ . In particular, we have that for all  $x \in \mathbb{Z}$ ,

$$
\{V(X)^n \delta_x : n \in \mathbb{N}\} \subseteq \text{span}\{\delta_x\}
$$

Clearly, this set cannot be dense in  $\mathcal{H}$ , and so no  $\delta_x$  is cyclic for  $V(X)$ .

(b) Note that for each  $x \in \mathbb{Z}$  and each  $z \in \mathbb{C}^+$  we have

$$
(V(X) - z\mathbb{1})\left(\frac{\delta_x}{V(x) - z}\right) = \delta_x \implies (V(X) - z\mathbb{1})^{-1}\delta_x = \left(\frac{1}{V(x) - z}\right)\delta_x
$$

So, we see that

$$
f_x(z) = \left\langle \delta_x, \left( \frac{1}{V(x) - z} \right) \delta_x \right\rangle = \frac{1}{V(x) - z}
$$

(c) From the above, we get that for all  $E \in \mathbb{R}$  and all  $\varepsilon > 0$ ,

$$
\operatorname{Im} \left\{ f_x(E + i\varepsilon) \right\} = \operatorname{Im} \left\{ \frac{1}{(V(x) - E) - i\varepsilon} \right\}
$$

Letting  $y := V(x) - E$  for notation, we can rationalize

$$
\frac{1}{y - i\varepsilon} = \frac{y + i\varepsilon}{y^2 + \varepsilon^2} = \frac{y}{y^2 + \varepsilon^2} + i\frac{\varepsilon}{y^2 + \varepsilon^2}
$$

From this, it is obvious that

$$
\operatorname{Im}\left\{f_x(E+i\varepsilon)\right\} = \frac{\varepsilon}{(V(x)-E)^2 + \varepsilon^2}
$$

Problem 9 continued on next page... 12

If  $E \neq V(x)$ , then we find that

$$
\lim_{\varepsilon \to 0^+} \text{Im} \left\{ f_x(E + i\varepsilon) \right\} = \lim_{\varepsilon \to 0^+} \varepsilon \text{Im} \left\{ f_x(E + i\varepsilon) \right\} = 0
$$

However, if  $E = V(x)$  then

$$
\lim_{\varepsilon \to 0^+} \mathbb{I}\mathrm{m} \left\{ f_x(E + i\varepsilon) \right\} = \lim_{\varepsilon \to 0^+} \frac{1}{\varepsilon} = \infty \quad \text{and} \quad \lim_{\varepsilon \to 0^+} \varepsilon \mathbb{I}\mathrm{m} \left\{ f_x(E + i\varepsilon) \right\} = 1
$$

(d) Let  $\mu_{V(X),\delta_0}$  be the spectral measure of  $(V(X),\delta_0)$ . Write the Radon-Nikodym decomposition of  $\mu_{V(X),\delta_0}$  w.r.t. the Lebesgue measure as

$$
\mu_{V(X),\delta_0} = f\lambda + \mu_{\text{sing}}
$$

By Lemma 10.10 in the lecture notes and part (c), we see that

$$
f(\lambda) = \lim_{\varepsilon \to 0^+} \frac{1}{\pi} \operatorname{Im} \{ f_0(\lambda + i\varepsilon) \} = \begin{cases} 0 & \lambda \neq V(0) \\ \infty & \lambda = V(0) \end{cases}
$$

So,  $f = 0$  Lebesgue a.e., which means  $f\lambda = 0$ . Therefore,  $\mu_{V(X),\delta_0}$  has no absolutely continuous part. By Lemma 10.11 in the lecture notes,  $spt(\mu_{sing}) = \{V(0)\}\$ , and at this value  $\mu_{V(X),\delta_0}$  has a point mass. From the above, we see that  $\mu_{sing}$  is precisely equal to a point mass at  $V(0)$ . Thus, the spectral measure  $\mu_{V(X),\delta_0}$  is pure point and equal to a point mass at  $V(0)$ .

On  $\mathcal{H} := \ell^2(\mathbb{N})$ , let R be the unilateral right shift operator. Calculate ker R, ker R<sup>\*</sup> and im(R) and show that  $R$  is a Fredholm operator. Calculate its Fredholm index.

### Solution

**Proof.** Let  $\{e_n\}_{n\in\mathbb{N}}$  be the standard orthonormal basis of H. Then, R is the operator sending

$$
Re_n := e_{n+1} \quad (n \in \mathbb{N})
$$

and extended linearly. Then, we see that R preserves the norm of each  $e_n$ , which means that it is an isometry over all of H by the Pythagorean theorem and the fact that  $\{e_n\}_n$  is an orthonormal basis. In particular, this means that  $R$  has trivial kernel and closed image. Let

$$
M := \text{span}\{e_1\} = \{\psi \in \mathcal{H} : \ \psi(j) = 0 \text{ for all } j > 1\}
$$

Then, we see that

$$
im(R) = M^{\perp} = \{ \psi \in \mathcal{H} : \psi(1) = 0 \}
$$

We have already seen that  $R^*$  is the unilateral left shift, i.e.

$$
R^*e_n := \begin{cases} e_{n-1} & n > 1 \\ 0 & n = 1 \end{cases}
$$

From this wee see immediately that  $\ker(R^*) = M$ .

Now, we see that  $\dim \ker(R) = 0$ . Also, we know that  $\dim \operatorname{coker}(R) = \dim \ker(R^*) = \dim M = 1$ . Therefore, *R* is Fredholm with a Fredholm index of  $0 - 1 = -1$ . ■

Show that on  $\mathcal{H} := \ell^2(\mathbb{N}), \frac{1}{X}$  where X is the position operator is not a Fredholm operator by calculating  $\text{im}(\frac{1}{X})$  and showing that it is not closed.

### Solution

**Proof.** Let  $\{e_n\}_{n\in\mathbb{N}}$  be the standard orthonormal basis of H. The operator  $\frac{1}{X}$  can be expressed as a multiplication operator on this basis via

$$
\frac{1}{X}e_n := \frac{1}{n}e_n
$$

Clearly,  $\frac{1}{X}$  is bounded and linear, and so  $\frac{1}{X} \in \mathcal{B}(\mathcal{H})$ . For any  $\psi \in \mathcal{H}$  we may write

$$
\psi=\sum_{n\in\mathbb{N}}\psi(n)e_n
$$

where this convergence is in the norm on  $H$ . So, we have that

$$
\frac{1}{X}\psi=\sum_{n\in\mathbb{N}}\frac{\psi(n)}{n}e_n
$$

Therefore, we may write

$$
\operatorname{im}\left(\frac{1}{X}\right) = \left\{\varphi \in \mathcal{H} : \sum_{n \in \mathbb{N}} n^2 |\varphi(n)|^2 < \infty\right\}
$$

since any such  $\varphi$  can be expressed as  $\frac{1}{X}\psi$  for a  $\psi \in \mathcal{H}$ , whereas any  $\varphi \in \mathcal{H}$  that does not have this summability condition cannot. We claim that this set is not closed. To see this, for each  $N \in \mathbb{N}$  define

$$
\psi_N:=\sum_{n=1}^N\frac{1}{n}e_n
$$

Clearly, each  $\psi_N \in \text{im}(\frac{1}{X})$  since it has finitely many nonzero terms. Furthermore, we see that  $\{\psi_N\}_N$  has a limit  $\psi$ , since

$$
\|\psi\|_{\mathcal{H}}^2=\Bigl\|\lim_{N\to\infty}\psi_N\Bigr\|_{\mathcal{H}}^2=\left\|\sum_{n=1}^\infty\frac{1}{n}e_n\right\|_{\mathcal{H}}^2=\sum_{n=1}^\infty\frac{1}{n^2}<\infty
$$

However, this limit is certainly not in  $\text{im}(\frac{1}{X})$  since  $\psi$  does not satisfy the harder summability condition

$$
\sum_{n \in \mathbb{N}} n^2 |\psi(n)|^2 = \sum_{n \in \mathbb{N}} n^2 \frac{1}{n^2} = \sum_{n \in \mathbb{N}} 1 = \infty
$$

So, there is a sequence of elements in  $\text{im}(\frac{1}{X})$  converging to a vector that is not in  $\text{im}(\frac{1}{X})$ , and so  $\text{im}(\frac{1}{X})$  is not closed. By Proposition 9.45 in the lecture notes, this means that  $\frac{1}{X}$  is not Fredholm. since  $coker(\frac{1}{X})$  is not finite-dimensional.  $\quad \blacksquare$